THE SIMULATION AND EXPERIMENTAL RESULTS OF DYNAMIC BEHAVIOUR OF TORQUE MOTOR HAVING PERMANENT MAGNETS

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The application, construction, and principle of operation of the torque motor having permanent magnets are presented first. Since dynamic behaviour of a servohydraulic system is, to a large extent, determined by dynamic behaviour of the applied torque motor, an analysis of motor's dynamic behaviour is carried out and its dynamic transfer characteristic is determined. Finally, the simulation and experimental results of dynamic behaviour of the torque motor are presented.

Key words: torque motor, permanent magnet, servo valve, dynamic model, simulation

1 INTRODUCTION

The most frequently applied electrically controlled hydraulic preamplifier is servovalve. Within servovalve an electric input signal is converted to a shift, or move, of a pilot piston, which is via a hydraulic amplifier converted to a hydraulic variable, usually flow or pressure. As a converter of the small current (or voltage) signal to mechanical movement, usually a controlled motor called torque motor is used.

Torque motor may produce a rotational or linear movement, where the developed torque, or force, is proportional to input current, [1].

In majority of servohydraulic valves the torque motor is an independent device. It is mounted separately and is hermetically sealed against the hydraulic part (operating fluid). The torque motor consists of the windings positioned across an anchor to provide its magnetization and permanent magnets which, together with polar extensions, constitute magnetic circuit of the armature.

By low power control signals (of the order 10 mV) one can control powers as large as several hundreds of kilowatts. A multitude of papers has been devoted to the control of hydraulic systems, their modeling, and improvement of their performance [1–16]. Dynamic behaviour of a system as a whole is largely, in addition to the hydraulic components, determined by dynamic processes in the torque motor, [5], [11–14]. In the present paper dynamic behaviour of torque motor is analyzed and its transfer characteristic in s -domain is determined. By the simulation and experimental results the conclusions of the dynamic analysis have been confirmed and the theoretical mathematical model of the torque motor has been verified.

2 DYNAMIC MODEL OF TORQUE MOTOR

The construction, principle of operation, and mathematical model of a torque motor having permanent magnets are given in [1]. If i – total signal current and if the armature windings are identical, current in each of the windings will be $i/2$. For each of the armature windings (Fig. 1) the following voltage relation applies

$$E_0 - R_0 i = R_c \frac{i}{2} + N_c \frac{d\Phi_a}{dt}, \quad (1)$$

where E_0 is the output voltage of the amplifier driving the torque motor, r_0 – amplifier's internal resistance, R_c – resistance of each of the windings, N_c – number of turns of each of the windings, and Φ_a – total magnetic flux passing through the armature.

Total magnetic flux Φ_a passing through the armature is given by the relation, derived elsewhere, which can be represented as

$$\Phi_a = 2\Phi_g \frac{x}{g} + \frac{N_c}{\mathfrak{R}_g} i \quad (2)$$

where: g – length of each of the air gaps when the armature is in neutral position, x – shift of armature ends with respect to the neutral position, $\mathfrak{R}_g = g/\mu_0 A_g$ – reluctance of each of the air gaps when the armature is in the neutral position, and Φ_g – flux in each of the air gaps when the armature is in the neutral position, [1].

From relations (1) and (2), by applying Laplace transform one obtains

$$E_0 = \left(R_0 + \frac{R_c}{2}\right)i(s) + K_b s\theta(s) + L_c s i(s). \quad (3)$$

K_b is voltage constant of each of the windings, expressed in (V/rad/s), given by relation

$$K_b = 2 \frac{a}{g} N_c \Phi_g. \quad (4)$$

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Fig. 1. Wiring diagram of torque motor

L_c is self-inductance of each of the windings, given by relation

$$L_c = \frac{N_c^2}{\mathfrak{R}_g}. \quad (5)$$

The relation for electromagnetic torque of work performed in the vicinity of operating point can be represented as [3]

$$T_d = K_t i + K_m \theta, \quad (6)$$

where: θ – angle of armature turning, and a – radius from the centre to armature ends, K_t is torque constant of the motor in (Nm/A), given by relation

$$K_t = 2 \frac{a}{g} N_c \Phi_g, \quad (7)$$

and K_m is a constant of magnetic elasticity of the motor in (Nm/rad), given by relation

$$K_m = 4 \left(\frac{a}{g} \right)^2 \mathfrak{R}_g \Phi_g^2. \quad (8)$$

By applying the second Newton law, one can write

$$T_d = J_a \frac{d^2 \theta}{dt^2} + B_a \frac{d\theta}{dt} + K_a \theta + T_L, \quad (9)$$

where: J_a – inertia of the armature and the attached moving parts, B_a – total viscous damping coefficient of the armature, load, and other moving masses, K_a – mechanical torsion spring constant of the armature pivot, and T_L – moment of load.

By combining (6) and (9) and by applying Laplace transform one obtains

$$K_t i(s) = J_a s^2 \theta(s) + B_a s \theta(s) + (K_a - K_m) \theta(s) + T_L. \quad (10)$$

Note that magnetic elasticity constant K_m has to be smaller than mechanical elasticity constant K_a in order that the system is stable. Relations (3) and (10), together with the torque motor parameters defined by (4), (5), (7), and (8) completely describe its dynamic behaviour. The

parameters are calculated on the basis of construction data for $\mathfrak{R}_g$, N_c , g , a , and ϕ_g . Variable ϕ_g is often expressed as product of flux density B_g in the air gap a surface A_g ($\phi_g = B_g A_g$). It is clear that for the defined parameters

$$K_t = K_b, \quad (11)$$

$$K_m L_c = K_b K_t. \quad (12)$$

By neglecting mechanical damping B_a , and by combining (3) and (10) one obtains

$$\theta(s) = \frac{K_o e_g(s) - (1 + \frac{s}{\omega_a}) T_L(s)}{(K_a (1 - \frac{K_m}{K_a})) (1 - \frac{K_m}{K_a})^{-1} \left(\frac{s^3}{\omega_a \omega_m^2} + \frac{s^2}{\omega_m^2} + \frac{s}{\omega_a} \right) + 1}, \quad (13)$$

where $K_o = \frac{K_t \mu}{(R_0 + \frac{R_c}{2}) K_a (1 - \frac{K_m}{K_a})}$ – static gain (in rad/V),

$\omega_a = \frac{2R_0 + R_c}{2L_c}$ – circuit-break frequency of the armature winding (rad/s),

$\omega_m = \sqrt{K_a/J_a}$ – mechanical natural frequency of the armature (rad/sec),

e_g – control voltage at input of the amplifier, and $\mu = E_0/e_g$ – voltage gain of the amplifier.

If numerical values of the coefficients are known, denominator of relation (13) can be factorized. The coefficients are so related that one can obtain a general form containing roots of the equation, which can be represented graphically. The final equation representing dynamic behaviour of the torque motor having control voltage e_g and load moment T_L is

$$\theta(s) = \frac{K_o e_g(s) - (1 + \frac{s}{\omega_a}) T_L(s)}{(1 + \frac{s}{\omega_r}) (\omega_o^2 + \frac{2\delta_o}{\omega_o} s + 1)}, \quad (14)$$

where ω_r , ω_o , and δ_o as functions of ω_a/ω_m are shown in diagrams of Figs. 2, 3, and 4, respectively. From these

Fig. 2. Normalized diagram of real roots of the torque motor transfer function

Fig. 3. Normalized curves of the natural frequency of the torque motor

Fig. 4. Plot of damping ratio of the torque motor

diagrams one can see that the third order characteristic equation has one linear and two subattenuating roots. These factors are the consequence of the electric time constant $1/\omega_a$ and mechanical resonance frequency of the armature. However, turning point frequencies of the third order characteristic equation are lower, *ie* $\omega_r < \omega_a$ and $\omega_o < \omega_m$. Since for realized torque motors $K_m/K_a < 1/2$, the following approximations are suitable for per-

forming the calculations

$$\omega_r \approx \omega_a, \quad \omega_o \approx \omega_m, \quad \delta_o \approx \frac{1}{2} \frac{K_m}{K_a}.$$

Transfer function of a load-free torque motor with respect to input voltage e_g is

$$\frac{\theta(s)}{e_g(s)} = K_o \left(1 + \frac{s}{\omega_r}\right)^{-1} \left(\frac{s^2}{\omega_o^2} + \frac{2\delta_o}{\omega_o} s + 1\right)^{-1}. \quad (15)$$

Usually this transfer function is given with respect to control current $i(s)$ in the form

$$\frac{\theta}{i} = \frac{K_t K_a^{-1} (1 - (K_m/K_a))^{-1}}{\omega_m^{-2} (1 - \frac{K_m}{K_a})^{-1} s^2 + B_a K_a^{-1} (1 - \frac{K_m}{K_a})^{-1} s + 1}. \quad (16)$$

One obtains second order transfer function having natural frequency slightly lower than ω_m . The attenuation is very small and depends upon structural attenuation of the torque spring.

3 SIMULATION MODEL

This section presents the simulation model of torque motor, which is applied to one type of servovalves of series 770 — product of MOOG. The model is built-in the program package PSPICE.

Equations (1), (2), (3), (9) and (10) are the basis for forming the model. Note that direct electromechanical analogies [4] have been used to represent mechanical elements by the electrical equivalents. On this basis the functional simulation scheme of the torque motor shown in Fig. 5 is given. Input amplifier is modeled by a voltage controlled voltage source E_o , of internal resistance R_o . The output variables are the accomplished torque, balanced by the load moment $T_L(t)$, and output deflection $\theta(t)$. The equivalent stiffness is given as difference between mechanical torque stiffness K_a and magnetic elasticity K_m .

4 THE SIMULATION RESULTS AND DISCUSSION

This section presents the characteristic waveforms of a specific torque motor which is used in servovalves of manufacturer MOOG — series 770. Specifications for this torque motor are taken from manufacturer's catalogue:

Electric data:

resistance of each winding at 25 °C $R_c = 22 \Omega$, self-inductance of winding $L_c = 0.06$ H, magnetic constant $K_m = 0.833$ A/Wb, torque constant, $K_t = 0.05$ Nm/rad, voltage constant $K_b = 0.05$ V/rad/s.

Mechanical data:

inertial torque of armature $J = 2 \times 10^{-6}$ kgm², elastic torsion spring constant $K_a = 1 \dots 4$ Nm/rad, and equivalent

Fig. 5. The simulation scheme of torque motor

Fig. 6. Characteristic waveforms for $K_m/K_a = 0.2$ Fig. 7. Characteristic waveforms for $K_m/K_a = 0.8$

Fig. 8. Bode diagrams of the torque motor

viscous damping of the armature $B_a = 0.0001$ N/m/s, radius of torque bar $a = 0.02$ m.

Figures 6 and 7 present the characteristic waveforms in the transient regime, while Fig. 8 shows the corresponding Bode diagrams, for the cases of two boundary ratios K_m/K_a . The attenuation grows with the growth of this ratio. However, the static linearity and hysteresis limit this ratio, thus this parameter has to be a compromise between the two requirements. In practice, this ratio is $K_m/K_a = 0.1 \dots 0.5$.

Figure 9 shows the characteristic waveforms for the case of sinusoidal current excitation of torque motor, when $K_m/K_a = 0.8$. The excitation frequency was tuned to $f_p = 20$ Hz.

Figure 10 shows the response of deflection of the torque motor, for a pulsed load moment $T_L = 2 \times 10^{-3}$ Nm, for the case $K_m/K_a = 0.2$, which is usual in real situations. The restoration time is about 100 ms, which represents a very good dynamic response.

Fig. 9. Characteristic waveforms for sinusoidal excitation; $f_p = 20$ Hz

Fig. 10. Response of the torque motor for a pulsed load moment

5 THE EXPERIMENTAL RESULTS

This Section presents a part of the experimental results related to the electric circuit for excitation of the torque motor, shown in Fig. 1. This circuit has been realized as control part of the hydraulic pulsation module applied within the system of the testing installation used for vibrational testing. The mentioned hydraulic servo-valves have been used in this system. The current of the torque motor and shift of the pilot piston have been measured for: sinusoidal excitation (Fig. 11) at excitation frequency $f_{pob} = 2$ Hz and square-wave excitation (Fig. 12) at excitation frequency $f_{pob} = 10$ Hz.

6 CONCLUSION

An analysis of dynamic behaviour of the torque motor used in servohydraulic systems has been carried out. The dynamic behaviour has to a large extent been influenced by the ratio of the mechanical torque and magnetic

Fig. 11. Oscilloscope traces at sinusoidal excitation

Fig. 12. Oscilloscope traces at square-wave excitation

constants. The obtained analytical results have been confirmed by the simulations and experiments.

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