

Numerical Solutions of the Modified Burgers' Equation by Finite Difference Methods

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Abstract

In this study, a numerical solution of the modified Burgers' equation is obtained by the finite difference methods. For the solution process, two linearization techniques have been applied to get over the non-linear term existing in the equation. Then, some comparisons have been made between the obtained results and those available in the literature. Furthermore, the error norms L_2 and L_∞ are computed and found to be sufficiently small and compatible with others in the literature. The stability analysis of the linearized finite difference equations obtained by two different linearization techniques has been separately conducted via Fourier stability analysis method.

Keywords: Modified Burgers' equation, Finite difference method, Fourier Stability Analysis.

AMS classification: 97N40, 65N30, 65D07, 76B25, 74S05, 74J35

1. INTRODUCTION

The one-dimensional generalized Burgers' equation is of the form

$$U_t + U^p U_x - \nu U_{xx} = 0, \quad a \leq x \leq b, t \geq 0$$

where $U(x, t)$ is the velocity for space x and time t , ν is a positive constant representing the kinematic viscosity of the fluid, and p is a positive parameter. When $p = 1$ we get Burgers' equation, $p = 2$ we get modified Burgers' equation.

The Burgers' equation has a wide range of applications in miscellaneous fields as a mathematical model for several phenomena and is thus of a great interest. The analytical and numerical solutions of the equation have been found out by several authors using various methods and techniques. In the present work, a variation of it has been considered, namely the modified Burgers' equation, given in the form of

$$U_t + U^2 U_x - \nu U_{xx} = 0, \quad a \leq x \leq b \quad (1)$$

where U is the dependent variable, ν is the viscosity parameter, and t and x are the independent parameters, denoting time and space, respectively. For the solution of the

numerical example, the following boundary conditions are going to be used

$$U(a, t) = \beta_1, \quad U(b, t) = \beta_2 \quad t \geq t_0. \quad (2)$$

The current work's main aim is to apply the finite difference methods to develop a numerical method for the approximate solution of the modified Burgers' equation. Eq. (1) has been solved both analytically and numerically by several authors using various methods and techniques. Some of them can be given as follows. The modified Burgers' equation has been solved by Ramadan and El-Danaf [2] using the collocation method with quintic splines. The equation has been numerically solved by Ramadan et al. [3] using the collocation method with septic splines. The Burgers' and modified Burgers' equations have been solved by Saka and Dag [4] by applying time and space splitting techniques and then employed the quintic B-spline collocation procedure to approximate the resulting systems. Irk [5] has employed Crank-Nicolson central differencing scheme for the time integration and sextic B-spline functions for the space integration to the modified and time splitted modified Burgers' equation. A numerical solution has been proposed by Temsah [6] for the convection-diffusion equation using El-Gendi method with interface points and then numerical results for Burgers' and modified Burgers' equations have been shown. Grienwank and El-Danaf [7] have proposed a non-polynomial spline based method to obtain numerical solutions of the non-linear modified Burgers' equation. Bratsos [8] has used a finite-difference scheme based on rational approximations to the matrix-exponential term in a two-time level recurrence relation for the numerical solution of the modified Burgers' equation. Bratsos [9] has presented a finite-difference scheme based on fourth-order rational approximants to the matrix-exponential term in a two-time level recurrence relation for the numerical solution of the modified Burgers equation. Bratsos and Petrakis [10] have used an explicit finite difference scheme based on second-order rational approximations to the matrix-exponential term for the numerical solution of the modified Burgers' equation. The equation has been numerically solved by Roshan and Bhamra [11] by the Petrov-Galerkin method using a linear hat function as the trial function and a cubic B-spline function as the test function.

In this study, two linearization techniques have been applied to deal with the non linear term while obtaining the numerical solution of the Modified Burgers' equation. A numerical example has been considered to test the performance of two

linearization techniques and then the stability analysis of the numerical schemes has been investigated separately.

2. THE FINITE DIFFERENCE METHOD

Let's suppose that the solution domain of the problem $a \leq x \leq b$ is divided into intervals having equal length h in the x direction and having equal time intervals k in time t such that $x_i = ih$, $i = 0(1)N$ and $t_j = jk$, $j = 0(1)J$ and U_{ij} will denote $U(x_i, t_j)$ throughout the article.

In the finite difference method, in place of the dependent variable and its derivatives their approximated values by the finite difference approximation are written. These approximations will result in either a single explicit equation or a system of difference equations. When applied to non-linear problems, it normally results in non-linear system of equations and they cannot be solved directly. Thus, an appropriate numerical algorithm is used to solve them.

3. LINEARIZATION I:

Using the forward difference approximation for U_t , the weighted central difference approximation for U_{xx} in Eq. (1) at the nodal point $(i, j+1)$

$$U_t \simeq \frac{U_{i,j+1} - U_{i,j}}{k},$$

and

$$U_{xx} \simeq \frac{1}{h^2} (\theta(U_{i+1,j+1} - 2U_{i,j+1} + U_{i-1,j+1}) + (1 - \theta)(U_{i+1,j} - 2U_{i,j} + U_{i-1,j})),$$

respectively, and applying the the following linearization technique for the non-linear term $U^2 U_x$

$$\begin{aligned} U^2 U_x \simeq & U_{i,j+1} U_{i,j} \left(\frac{U_{i+1,j} - U_{i-1,j}}{2h} \right) + U_{i,j} U_{i,j+1} \left(\frac{U_{i+1,j} - U_{i-1,j}}{2h} \right) + \\ & U_{i,j} U_{i,j} \left(\frac{U_{i+1,j+1} - U_{i-1,j+1}}{2h} \right) - 2U_{i,j} U_{i,j} \left(\frac{U_{i+1,j} - U_{i-1,j}}{2h} \right), \end{aligned}$$

we can easily obtain the following system of algebraic equations

$$\begin{aligned} & \frac{U_{i,j+1} - U_{i,j}}{k} + U_{i,j+1}U_{i,j}\left(\frac{U_{i+1,j} - U_{i-1,j}}{2h}\right) + U_{i,j}U_{i,j+1}\left(\frac{U_{i+1,j} - U_{i-1,j}}{2h}\right) + \\ & U_{i,j}U_{i,j}\left(\frac{U_{i+1,j+1} - U_{i-1,j+1}}{2h}\right) - 2U_{i,j}U_{i,j}\left(\frac{U_{i+1,j} - U_{i-1,j}}{2h}\right) - \\ & \frac{\nu}{h^2}(\theta(U_{i+1,j+1} - 2U_{i,j+1} + U_{i-1,j+1}) + (1-\theta)(U_{i+1,j} - 2U_{i,j} + U_{i-1,j})) = 0, \end{aligned} \quad (3)$$

for $i = 1(1)N - 1$ and $j = 0(1)J$.

For different values of θ ($\theta = 0, 1/2, 1$), the Eq. (3) is going to be solved using an appropriate algorithm.

3.1. Stability analysis

To investigate the stability of the approximation obtained by the present algorithm, we will use the von Neumann theory in which the growth factor of a typical Fourier mode is defined as:

$$U_m^n = e^{i\beta^{ph}\xi^q}, \quad (4)$$

where $i = \sqrt{-1}$. To investigate the stability of the numerical scheme, the nonlinear term U^2U_x in the modified Burgers' equation has been linearized by making the quantity U^2 a local constant. Thus the nonlinear term in the equation converts into $\widehat{U}U_x$ and the Eq. (1) becomes

$$U_t + \widehat{U}U_x - \nu U_{xx} = 0$$

If we take the weighted average approximation as

$$\begin{aligned} & \frac{U_{m,n+1} - U_{m,n}}{k} + \widehat{U}\left(\theta\left(\frac{U_{m+1,n+1} - U_{m-1,n+1}}{2h}\right) + \right. \\ & \left. + (1-\theta)\left(\frac{U_{m+1,n} - U_{m-1,n}}{2h}\right)\right) - \frac{\nu}{h^2}(\theta(U_{m+1,n+1} - 2U_{m,n+1} + U_{m-1,n+1}) + \\ & \left. + (1-\theta)(U_{m+1,n} - 2U_{m,n} + U_{m-1,n})) = 0, \end{aligned} \quad (5)$$

the generalized m^{th} row of Eq. (5) becomes

$$\begin{aligned} & U_{m-1}^{n+1}\left(-\frac{\theta\widehat{U}}{2h} - \frac{\nu\theta}{h^2}\right) + U_m^{n+1}\left(\frac{1}{k} + \frac{2\nu\theta}{h^2}\right) + U_{m+1}^{n+1}\left(\frac{\theta\widehat{U}}{2h} - \frac{\nu\theta}{h^2}\right) \\ & = U_{m-1}^n\left(\frac{(1-\theta)\widehat{U}}{2h} + \frac{\nu(1-\theta)}{h^2}\right) + U_m^n\left(\frac{1}{k} - \frac{2\nu(1-\theta)}{h^2}\right) + \\ & + U_{m+1}^n\left(-\frac{(1-\theta)\widehat{U}}{2h} + \frac{\nu(1-\theta)}{h^2}\right). \end{aligned} \quad (6)$$

Substituting the Fourier mode (4) into the linearised recurrence relationship (3)

yields

$$g = \frac{a - ib}{c - id} \quad (7)$$

where

$$\begin{aligned} a &= h^2 - 2kv + 2kv\theta - 2kv(\theta - 1)\cos\phi, \\ b &= -kh\widehat{U}(\theta - 1)\sin\phi, \\ c &= h^2 + 2kv\theta - 2kv\theta\cos\phi, \\ d &= -kh\widehat{U}\theta\sin\phi. \end{aligned} \quad (8)$$

If $\theta = 0$ is taken, it corresponds to explicit method, and the following inequality is required for the stability condition.

$$h^4 - (h^2 - 2kv + 2kv\cos\phi)^2 - h^2k\widehat{U}^2\sin^2\phi \geq 0$$

If $\theta = 1$ is taken, it corresponds to implicit method, and the following inequality is required for the system to be stable.

$$4h^2kv + 4k^2v^2 - 4h^2kv\cos\phi - 8k^2v^2\cos\phi + 4k^2v^2\cos^2\phi + h^2k^2\widehat{U}^2\sin^2\phi \geq 0$$

If $\theta = \frac{1}{2}$ is taken, it corresponds to Crank-Nicolson method, and the scheme is unconditionally stable by the following inequality

$$-4h^2kv(\cos\phi - 1) \geq 0$$

After some basic arithmetic operations, it is seen that the stability condition $|g| \leq 1$ is satisfied by the following inequality:

$$c^2 + d^2 - a^2 - b^2 = 96h^2v\Delta t(2 + \cos\phi)\sin\left[\frac{\phi}{2}\right]^2 \geq 0$$

therefore we have come to the conclusion that the linearised scheme is unconditionally stable.

4. LINEARIZATION II:

Eq. (1) can be written as

$$\frac{\partial U}{\partial t} + \frac{1}{p+1} \frac{\partial U^{p+1}}{\partial x} - v \frac{\partial^2 U}{\partial x^2} = 0.$$

Using the forward difference approximation for U_t , the Crank-Nicolson difference approximations for $(U^{p+1})_x$ and U_{xx} , and then utilizing the central difference operator δ defined by $\delta_x U_{m,n} = U_{m+1,n} - U_{m-1,n}$ (see, e.g. [14]), Eq. (1) yields the system of algebraic equations

$$\begin{aligned} & \frac{U_{m,n+1} - U_{m,n}}{k} + \frac{1}{4h(p+1)} \left\{ \delta_x \left(U_{m,n+1}^{p+1} \right) + \delta_x \left(U_{m,n}^{p+1} \right) \right\} - \\ & - \frac{v}{2h^2} (U_{m+1,n+1} - 2U_{m,n+1} + U_{m-1,n+1} + U_{m+1,n} - 2U_{m,n} + U_{m-1,n}) = 0. \end{aligned} \quad (9)$$

for $m = 1(1)M-1$ and $n = 0(1)N$.

4.1. Stability analysis

To investigate the stability of the above scheme, we perform the computation of Eq. (9) with the values $U_{m,n}^*$ instead of $U_{m,n}$. Introducing an error $E_{m,n}$ given by $E_{m,n} = U_{m,n}^* - U_{m,n}$ and substituting it into Eq. (9) leads to

$$\begin{aligned} & U_{m,n+1}^* - U_{m,n+1} - (U_{m,n}^* - U_{m,n}) + \frac{k}{(p+1)4h} \left\{ U_{m+1,n+1}^{*p+1} - U_{m+1,n+1}^{p+1} - \right. \\ & - \left(U_{m-1,n+1}^{*p+1} - U_{m-1,n+1}^{p+1} \right) + U_{m+1,n}^{*p+1} - U_{m+1,n}^{p+1} - \left(U_{m-1,n}^{*p+1} - U_{m-1,n}^{p+1} \right) \left. \right\} - \\ & - \frac{vk}{2h^2} \left\{ U_{m+1,n+1}^* - U_{m+1,n+1} - 2(U_{m,n+1}^* - U_{m,n+1}) + \right. \\ & + U_{m-1,n+1}^* - U_{m-1,n+1} + U_{m+1,n}^* - U_{m+1,n} - 2(U_{m,n}^* - U_{m,n}) + U_{m-1,n}^* - U_{m-1,n} \left. \right\} = 0. \end{aligned}$$

We now assume that U varies little over a small region in comparison with the errors means that

$$U_{m-1,n} \simeq U_{m,n} \simeq U_{m+1,n} \simeq U_{m-1,n+1} \simeq U_{m,n+1} \simeq U_{m+1,n+1}.$$

It is also assumed that $E_{m,n}$ is sufficiently small compared with $U_{m,n}$ and then

$$U_{m,n}^{*p+1} - U_{m,n}^{p+1} = (U_{m,n} + E_{m,n})^{p+1} - U_{m,n}^{p+1} \simeq (p+1)E_{m,n}U_{m,n}^p$$

for all m and n . Using the above assumptions and substituting Fourier mode $E_{m,n} =$

$\xi^n e^{i\beta m h}$, ($i = \sqrt{-1}$) into the scheme gives the growth factor ξ of the form

$$\xi = \frac{1 - 4A \sin^2\left(\frac{\beta h}{2}\right) - i2B \sin(\beta h)}{1 + 4A \sin^2\left(\frac{\beta h}{2}\right) + i2B \sin(\beta h)} \quad (10)$$

where $A = \frac{vk}{2h^2}$ and $B = \frac{kU_{m,n}^p}{4h}$. Taking the modulus of (10) gives $|\xi| \leq 1$. The scheme is therefore unconditionally stable.

Clearly, the scheme (9) is a non-linear system of equations in $U_{m,n+1}$ and it needs to use an iterative technique to evaluate the solution. The main aim of this study is to solve the scheme (9) by a direct method. Using a Taylor series expansion of $U_{m,n+1}^{p+1}$ about the point (m, n) we obtain

$$\begin{aligned} U_{m,n+1}^{p+1} &= U_{m,n}^{p+1} + k \frac{\partial U_{m,n}^{p+1}}{\partial t} + \dots \\ &= U_{m,n}^{p+1} + k \frac{\partial U_{m,n}^{p+1}}{\partial U_{m,n}} \frac{\partial U_{m,n}}{\partial t} + \dots \end{aligned}$$

Hence in terms of order k , $U_{m,n+1}^{p+1} \cong U_{m,n}^{p+1} + (p+1)U_{m,n}^p (U_{m,n+1} - U_{m,n})$ and taking

$$W_m = U_{m,n+1} - U_{m,n} \quad (11)$$

Eq. (9), with some manipulations, leads to

$$\begin{aligned} \left(\frac{1}{4h}U_{m-1,n}^p + \frac{v}{2h^2}\right)W_{m-1} - \left(\frac{1}{k} + \frac{v}{h^2}\right)W_m + \left(\frac{v}{h^2} - \frac{1}{4h}U_{m+1,n}^p\right)W_{m+1} = \\ \frac{1}{2h(p+1)}\left(U_{m+1,n}^{p+1} - U_{m-1,n}^{p+1}\right) - \frac{v}{h^2}(U_{m+1,n} - 2U_{m,n} + U_{m-1,n}), \end{aligned} \quad (12)$$

$(m = 1(1)M - 1)$ a system of linear equations for W_m . This approximation is second order in both space and time as regards truncation error. Obviously, the solution at the $(n+1)th$ time level is obtained from (11) as $U_{m,n+1} = U_{m,n} + W_m$ [14].

5. NUMERICAL EXAMPLES AND RESULTS

For the test problem used in the present work, numerical results of the equation have been obtained and all computations have been run on a Pentium i7 PC in the Fortran code using double precision arithmetic. To show how accurate the results, both the error norm L_2

$$L_2 = \|U^{exact} - U_N\|_2 = \sqrt{h \sum_{j=0}^N |U_j^{exact} - (U_N)_j|^2}, \quad (13)$$

and the error norm L_∞

$$L_\infty = \|U^{exact} - U_N\|_\infty = \max_j |U_j^{exact} - (U_N)_j|. \quad (14)$$

are going to be computed and presented.

6. TEST PROBLEM

The analytical solution of the modified Burgers' equation is given as

$$U(x, t) = \frac{x/t}{1 + \sqrt{t}/c_0 \exp(x^2/4vt)}, \quad t \geq t_0, \quad 0 \leq x \leq 1 \quad (15)$$

where c_0 is a constant, $0 < c_0 < 1$ and $t_0 = 1$.

For the initial condition of test problem, we will take equation (15) by evaluating it at $t = 1$. For the boundary conditions, we will use $U(0, t) = U_x(0, t) = 0$ and $U(1, t) = U_x(1, t) = 0$. Various viscosity constants $v = 0.01, 0.001, 0.005$, space steps $h = 0.005$, time steps $\Delta t = 0.01$ and $c_0 = 0.5$ will be taken over the problem domain $[0, 1]$ during the solution process of the problem. First of all, the program has been run until the time $t = 11$ and then the error norms L_2 and L_∞ are computed and presented in Table I for different values of viscosity v . As it is seen from the table, both of the error norms L_2 and L_∞ are small enough. Both of the error norms L_2 and L_∞ have been compared with those of some other authors for various values of h and v in Table II. It is clearly seen from the table that both of the error norms are better or as good as the others found in literature.

Table I. Comparison of the error norms L_2 and L_∞ with $h = 0.005$ and $\Delta t = 0.01$ for various values of ν .

t	$\nu = 0.01$		$\nu = 0.005$		$\nu = 0.001$	
	$L_2 \times 10^3$	$L_\infty \times 10^3$	$L_2 \times 10^3$	$L_\infty \times 10^3$	$L_2 \times 10^3$	$L_\infty \times 10^3$
1	0.000000	0.000000	0.000000	0.000000	0.000000	0.000000
2	0.378848	0.816262	0.225949	0.579150	0.067286	0.259046
3	0.344560	0.709949	0.205429	0.503503	0.061409	0.225358
4	0.317165	0.605190	0.188055	0.429111	0.056346	0.192213
5	0.307896	0.526341	0.175034	0.372673	0.052504	0.166926
6	0.326003	0.525791	0.164589	0.329766	0.049394	0.147826
7	0.369938	0.755043	0.155888	0.296209	0.046765	0.132819
8	0.427983	0.963399	0.148688	0.269279	0.044486	0.120756
9	0.489147	1.139612	0.143137	0.247219	0.042478	0.110820
10	0.547020	1.281253	0.139607	0.228838	0.040688	0.102585
11	0.598717	1.390450	0.138473	0.213415	0.039078	0.095516

Table II. Comparison of the error norms L_2 and L_∞ with those in other studies in the literature at $t = 2, 6, 10$.

	$t = 2$		$t = 6$		$t = 10$	
	$L_2 \times 10^3$	$L_\infty \times 10^3$	$L_2 \times 10^3$	$L_\infty \times 10^3$	$L_2 \times 10^3$	$L_\infty \times 10^3$
$h=0.005, \Delta t = 0.01, \nu = 0.01$						
EFD($\Delta t = 0.001$)	0.37869	0.81610	0.32600	0.52579	0.54702	1.28125
IFD	0.39169	0.83158	0.32683	0.52579	0.54712	1.28125
CN-FD	0.37986	0.81754	0.32608	0.52579	0.54704	1.28125
LFD	0.37978	0.81736	0.32605	0.52579	0.54702	1.28125
[2]	0.52308	1.21698	0.49023	0.72249	0.64007	1.28124
[3]	0.79043	1.70309	0.57672	0.76105	0.80026	1.80329
[5], (SBCM1)	0.38489	0.82934	-	-	0.54826	1.28127
[5], (SBCM2)	0.39078	0.82734	-	-	0.54612	1.28127
EFD($\Delta t = 0.001$), [0,1,3]	0.37869	0.81610	0.27624	0.46512	0.25308	0.32449
IFD, [0,1,3]	0.39169	0.83158	0.27666	0.46734	0.25391	0.32500
CN-FD, [0,1,3]	0.37986	0.81754	0.27630	0.46533	0.25397	0.32455
LFD, [0,1,3]	0.37978	0.81735	0.27627	0.46528	0.25395	0.32452
[5], (SBCM1), [0,1,3]	0.38489	0.82934	-	-	0.25586	0.32723
[5], (SBCM2), [0,1,3]	0.39078	0.82734	-	-	0.25259	0.32337
$h=0.005, \Delta t = 0.001, \nu = 0.005$						
EFD	0.22643	0.57988	0.16460	0.32987	0.13960	0.22886
IFD	0.22781	0.58193	0.16464	0.33014	0.13958	0.22891
CN-FD	0.22712	0.58091	0.16462	0.33000	0.13959	0.22888
LFD	0.22711	0.58090	0.16462	0.33000	0.13959	0.22888
[2]	0.25786	0.72264	0.22569	0.43082	0.18735	0.30006
[5], (SBCM1)	0.22890	0.58623	-	-	0.14042	0.23019
[5], (SBCM2)	0.23397	0.58424	-	-	0.13747	0.22626
$h=0.005, \Delta t = 0.01, \nu = 0.001$						
EFD	0.06695	0.25830	0.04939	0.14773	0.04070	0.10257
IFD	0.07114	0.26743	0.04951	0.14895	0.04061	0.10277
CN-FD	0.06900	0.26287	0.04944	0.14834	0.04065	0.10267
LFD	0.06900	0.26284	0.04944	0.14833	0.04065	0.10267
[2]	0.06703	0.27967	0.06046	0.17176	0.05010	0.12129
[3]	0.18355	0.81862	0.08142	0.21348	0.05512	0.13943
[5], (SBCM1)	0.06843	0.26233	-	-	0.04080	0.10295
[5], (SBCM2)	0.07220	0.25975	-	-	0.03871	0.09882
$h=0.02, \Delta t = 0.01, \nu = 0.01$						
EFD	0.37559	0.80866	0.32916	0.52579	0.55844	1.28125
IFD	0.39938	0.83959	0.33067	0.52579	0.55861	1.28125
CN-FD	0.38724	0.82351	0.32988	0.52579	0.55852	1.28125
LFD	0.38717	0.82328	0.32982	0.52579	0.55849	1.28125
[3]	0.79043	1.70309	0.51672	0.76105	0.80026	1.80239
[5], (SBCM1)	0.38474	0.82611	-	-	0.55985	1.28127
[5], (SBCM2)	0.41321	0.81502	-	-	0.55095	1.28127

The computed numerical results together with their errors are graphed in Figures 1-3 for various values of ν at different time levels. But the graphs of the errors have only been drawn at time $t = 10$. It can be seen that the maximum error happens at the right-hand boundary of the solution domain for $\nu = 0.01$. However, the errors for $\nu = 0.005$ and $\nu = 0.001$ have been recorded around the the points where the waves

get their highest amplitudes.

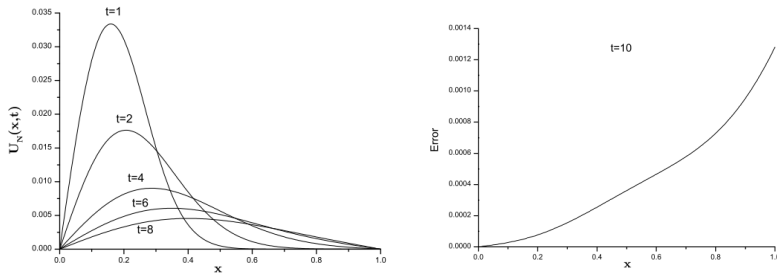


Fig. 1. The numerical solutions of Problem at different times with $\nu = 0.01$.

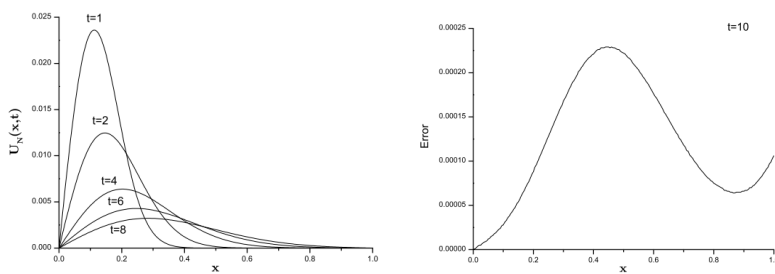


Fig. 2. The numerical solutions of Problem at different times with $\nu = 0.005$.

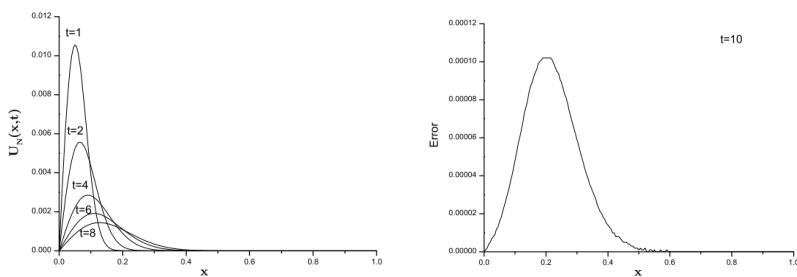


Fig. 3. The numerical solutions of Problem at different times with $\nu = 0.001$.

7. CONCLUSIONS

In the present work, numerical solutions of the Modified Burgers' equation based on the finite difference methods have been presented. To show the performance of the presented algorithm, a test problem has been considered. The performance and efficiency of the method are shown by calculating the error norms L_2 and L_∞ . The obtained results show that the error norms are sufficiently small during all computer runs. The obtained results indicate that the present method is a particularly successful numerical scheme to solve the Modified Burgers' equation. Therefore, the method used in the present work can strongly be advised to get approximate solutions of several other widely used non-linear equations in the literature.

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